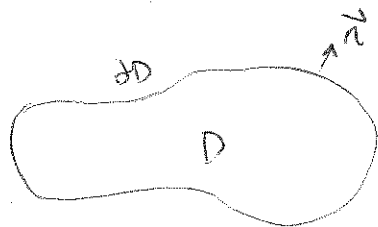


## §7.1 Green's First Identity

The divergence theorem:

$$\iiint_D \operatorname{div} \vec{F} \, d\vec{x} = \iint_{\partial D} \vec{F} \cdot \vec{n} \, dS$$

where  $\vec{F}$  is any vector function,  $D$  is a bounded solid region and  $\vec{n}$  is the unit outer normal on  $\partial D$ .



Starting from the product rule,

$$(v u_x)_x = v_x u_x + v u_{xx}$$

we sum  $x, y, z$  derivatives

$$(v u_x)_x + (v u_y)_y + (v u_z)_z = v_x u_x + v_y u_y + v_z u_z + v(u_{xx} + u_{yy} + u_{zz})$$

In vector form,  $\nabla \cdot (v \nabla u) = \nabla v \cdot \nabla u + v \Delta u$ .

We integrate and use the divergence theorem on the left

side to get 
$$\iint_{\partial D} v \frac{\partial u}{\partial n} \, dS = \iiint_D \nabla v \cdot \nabla u \, d\vec{x} + \iiint_D v \Delta u \, d\vec{x}$$

Green's first identity (G1)

If  $v \equiv 1$ , we have 
$$\iint_{\partial D} \frac{\partial u}{\partial n} \, dS = \iiint_D \Delta u \, d\vec{x}$$

Consider the Neumann problem

$$\begin{cases} \Delta u = f(x) & \text{in } D \\ \frac{\partial u}{\partial n} = h(x) & \text{on } \partial D \end{cases} \quad \iint_{\partial D} h \, dS = \iiint_D f \, d\vec{x}$$

← compatibility condition

Mean value property: the average value of any harmonic function over a sphere equals its value at the center.

Proof: Without loss of generality,  $D = \{ |\vec{x}| < a \}$   $\partial D = \{ |\vec{x}| = a \}$

$$\vec{n} = \frac{\vec{x}}{|\vec{x}|} = \frac{\vec{x}}{r} \quad \frac{\partial u}{\partial n} = \nabla u \cdot \vec{n} = \frac{\partial u}{\partial x} \frac{x}{r} + \frac{\partial u}{\partial y} \frac{y}{r} + \frac{\partial u}{\partial z} \frac{z}{r} = \frac{\partial u}{\partial r}$$

(1)

$$\iint_{\partial D} \frac{\partial u}{\partial r} dS = 0$$

In spherical coordinates  $\int_0^{2\pi} \int_0^\pi u(r, \theta, \phi) a^2 \sin \theta d\theta d\phi = 0$   
 Since the area of  $\partial D$  is  $4\pi a^2$ ,

$$\frac{1}{4\pi a^2} \int_0^{2\pi} \int_0^\pi u(r, \theta, \phi) a^2 \sin \theta d\theta d\phi = 0$$

$$\rightarrow \frac{d}{dr} \left[ \frac{1}{4\pi} \int_0^{2\pi} \int_0^\pi u(r, \theta, \phi) \sin \theta d\theta d\phi \right] \Big|_{r=a} = 0 \quad \forall a > 0$$

so  $\frac{1}{4\pi} \int_0^{2\pi} \int_0^\pi u(r, \theta, \phi) \sin \theta d\theta d\phi$  is independent of  $r$ .

Let  $r \rightarrow 0$ , we get

$$\frac{1}{4\pi} \int_0^{2\pi} \int_0^\pi u(\vec{0}) \sin \theta d\theta d\phi = u(\vec{0})$$

That is

$$\begin{aligned} u(\vec{0}) &= \frac{1}{4\pi} \int_0^{2\pi} \int_0^\pi u(r, \theta, \phi) \sin \theta d\theta d\phi \\ &= \frac{1}{|\partial D|} \iint_{\partial D} u dS \end{aligned}$$

Remark: The maximum principle can be proved by the mean value property. Moreover,  $\frac{\partial u}{\partial n} > 0$  at the maximum point, which is called Hopf maximum principle.

Uniqueness of Dirichlet's Problem by Energy Method.

Assume two harmonic functions with the same boundary data.

Define  $u = u_1 - u_2$  harmonic and zero boundary data.

From (6.1), choosing  $v = u$ , we have

$$0 = \iint_{\partial D} u \frac{\partial u}{\partial n} dS = \iiint_D |\nabla u|^2 dx$$

By the first vanishing theorem, it follows that  $|\nabla u|^2 \equiv 0$

in  $D$ , then  $u \equiv C$  in  $D$ . Since  $u = 0$  on  $\partial D$ ,

$u \equiv 0$  in  $D$ . This proves the uniqueness.

Uniqueness of Neumann's Problem: If  $\Delta u = 0$  in  $D$  and  $\frac{\partial u}{\partial n} = 0$  on  $\partial D$ , then  $u$  is a constant in  $D$ .

Dirichlet's Principle:

Define the energy as  $E[w] = \frac{1}{2} \iiint_D |\nabla w|^2 d\vec{x}$

where  $w(x)$  defined in  $D$  satisfies the Dirichlet BC

$$w = h(\vec{x}) \text{ on } \partial D.$$

Let  $u(\vec{x})$  be the unique harmonic function in  $D$  that satisfies  $u = h(\vec{x})$  on  $\partial D$ . Let  $w(\vec{x})$  be any function that satisfies  $w = h(\vec{x})$  on  $\partial D$ . Then

$$E[w] \geq E[u]$$

$u$  is called ground state in some physical contexts and  $E[u]$  is the ground state energy.

Proof: Let  $v = u - w$ , then

$$E[w] = \frac{1}{2} \iiint_D |\nabla(u-w)|^2 d\vec{x}$$

$$= E[u] - \iiint_D \nabla u \cdot \nabla v d\vec{x} + E[v]$$

Since  $\Delta u = 0$  in  $D$   $v = 0$  on  $\partial D$ , by GL, we have

$$E[w] = E[u] + E[v] \geq E[u] \quad \#$$